

COMMUNICATION

Full Solution Processed Mesostructured Optical Resonators Integrating Colloidal Semiconductor Quantum Dots

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Mauricio E. Calvo,^a Nuria Hidalgo,^a Roland Schierholz,^{a,b} András Kovács,^c Asunción Fernández,^a Martín G. Bellino,^d Galo J.A.A. Soler-Illia,^d Hernán Míguez,^{a*}

Herein we show a solution based synthetic pathway to obtain a resonant optical cavity with embedded colloidal semiconductor quantum dots (CSQDs). The optical cavity pore network, surrounded by two dense Bragg mirrors, was designed *ad hoc* to selectively host the quantum dots, while uncontrolled infiltration of those in the rest of the layered structure was prevented. Coupling between the optical resonant modes of the host and the natural emission of the embedded nanoparticles gives rise to the fine tuning of the luminescence spectrum extracted from the ensemble. Our approach overcomes, without the need of an encapsulating agent and exclusively by solution processing, the difficulties that arise from the low thermal and chemical stability of the CSQDs. It opens the route to achieving precise control over their location and hence over the spectral properties of light emitted by these widely employed nanomaterials. Furthermore, as the porosity of the cavity is preserved after infiltration, the system remains responsive to environmental changes, which constitutes an added value of the proposed structure.



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Introduction

Colloidal semiconductor quantum dots (CSQDs) present intense photoemission that can be precisely tuned by means of particle size and shape, as well as by modification of their capping layer.^{1,2,3} Another means of tailoring (spectrally or directionally) the photoluminescence of these nanomaterials is by engineering their photonic environment. In this regard, several works report on the modification of the luminescence of CSQDs infiltrated in porous 2D or 3D photonic crystals (PCs). In the case of 2D PCs, the effect of cavity modes on both emission intensity and lifetime have been investigated by integrating CSQDs either in the bulk of the structure^{4,5,6} or in built-in point-like and line defects.^{7,8} A variant of this approach makes use of the integration of CSQDs in 1D PC waveguides.⁹ In all these cases, sophisticated microfabrication techniques are employed to prepare the luminescent resonators (radio frequency and magnetron sputtering, electron beam lithography, molecular beam epitaxy), and combined with solution processing of CSQDs (in solvents, liquid crystals or polymers). These approaches significantly limit the integration of luminescent optical cavities based on quantum dots over large areas and onto arbitrary composition substrates, hence restricting the impact of the micrometre size structures developed to date to mere proofs of concept. In the case of 3D self-assembled PCs, CSQDs are in general homogeneously distributed throughout the structure and luminescent properties are the results of the coupling to band edges, low dispersion modes or band gaps at high energies.¹⁰⁻¹² Structures present a high density of intrinsic defects that smears the quality and magnitude of the targeted effects.

Interestingly, and although it would be a simpler alternative way to attain large area coatings with tailored photoemission, as it has been demonstrated for other luminescent species,¹³⁻¹⁵ very few works address the inclusion of CSQDs in planar 1D PC cavities, i.e., a slab surrounded by two multi-layered mirrors.^{16,17} The reason for this is the low compatibility between the commonly employed preparation methods to build a Bragg stack of high dielectric contrast thin films and the usually low thermal and chemical stability of CSQDs. In general, a multilayer structure is built by sequential deposition of the constituent films from either gas, by physical or chemical deposition methods, or liquid phase, by approaches based on sol-gel techniques. Typically, CSQDs do not withstand the chemical reactions occurring during deposition or the later thermal or chemical stabilization of the slabs. Hence, apart from the abovementioned cases, optical cavity multilayers made of liquid crystal films have been used as hosts for CSQDs,¹⁸ this being the only example of full liquid processing of an optical resonator containing CSQDs. Related to this, thin layers of CSQDs have been deposited onto multilayers.¹⁹

Herein we demonstrate a solution based synthetic route to prepare optical resonators with CdSe/ZnS core-shell colloidal QDs embedded at the desired depth. We make use of recent progress

in the field of supramolecularly templated mesoporous layers²⁰ to create a highly accessible optical cavity, sandwiched between two dense optical multilayers, with pores large enough as to selectively integrate organically capped CdSe/ZnS nanocrystals *a posteriori*, hence preventing their chemical or thermal degradation. Although inclusion of CSQDs within mesoporous layers can be achieved by *in situ* preparation through adsorption-reaction protocols,²¹ this route usually leads to long wavelength emission, irrespective of nanoparticle size, due to surface defects.²² Our approach overcomes this obstacle by directly incorporating pre-formed nanoparticles with well-defined luminescence properties (i.e., band position and intensity). Infiltration of these CSQDs is achieved by soaking a highly accessible resonator in the corresponding nanocrystal suspension. In this resonator, only the middle defect layer, which acts as optical cavity, presents an open interstitial structure and its pores are designed to host relatively large particles (maximum diameter ≈ 10 nm). This improves the selectivity and simplifies the analysis of the optical response caused by variations in the environment with respect to structures in which the CSQDs are homogeneously infiltrated in the whole photonic structure.^{11,12,15} The photoemission of CdSe/ZnS nanocrystals integrated in the resonator is strongly modified by the interplay with the cavity modes. A direct correlation is observed between the intensity and shape of the luminescence spectra measured at different collection directions and the angular dependence of the cavity resonant mode. Furthermore, as the middle layer remains open after CSQD infiltration, the resonator responds to changes in the environment with a different luminescent emission, as we herein demonstrate. Our approach opens the route to tailoring light emission of CSQDs by using several square centimetre large nanostructured coatings compatible with current sensing and optoelectronics technology.²³

Results and discussion

The scheme illustrating both the sequential procedure and intermediate layered materials realized to prepare the resonator and its nanostructure is depicted in Figure 1. Periodic layered structures were attained by dip-coating different substrates with silicon dioxide (SiO₂) and titanium dioxide (TiO₂) precursors. By this means, we obtain a dense one dimensional photonic crystal,²⁴⁻²⁶ which has been shown to be compatible with the integration of mesostructured films to yield more complex structures with devised optical properties.^{27,28} In this case, the deposition of an intercalated mesostructured film was realized also by dip-coating a hydrolyzed TiO₂ precursor mixed with a supramolecular template.^{29,30} A combination of surfactants, Pluronic F127 and poly(propylene glycol) (PPG), in a THF-butanol solution was chosen. It has been previously proven that this method permits to attain a bimodal, low order, interconnected mesopore network once the organic mold is eliminated.³¹ A detailed description of the preparation methods are provided in



the experimental section. We chose this material to form the optical cavity because its reported pore network allowed us foreseeing a suitable diffusion of small particles through it. From the perspective of the optical design, other middle layer compositions would also give rise to effects similar to those herein reported. Quantum dots were arbitrarily selected to emit in the middle of the visible range.

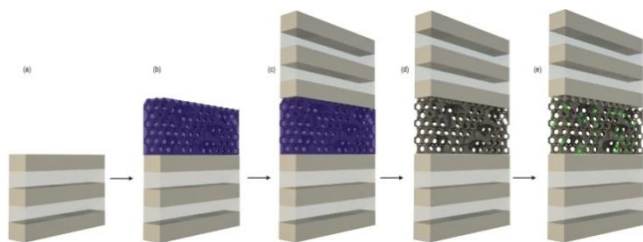


Fig. 1 Scheme of the process to integrate CSQD in a mesoporous optical resonator. (a) initial multilayer structure, (b) deposition of a bimodal mesostructured films, (c) deposition of the upper multilayer, (d) template removal, and (e) colloidal quantum dot infiltration

Chemical element sensitive annular dark field (ADF) scanning transmission electron microscopy (STEM) and high resolution TEM (HRTEM) images of the cross section of the multilayer made by sandwiching a mesostructured TiO_2 layer in between two 3-unit cell periodic dense SiO_2 - TiO_2 layers and a detail of the porous optical cavity are shown in Figures 2.a and 2.b, respectively. The contrast between the electronic density of both types of layers allows to clearly identify them as bright (TiO_2) and dark (SiO_2) fringes in Figure 2.a. Regarding their crystalline phase, X-ray diffraction (XRD) showed that SiO_2 layers were vitreous while TiO_2 ones presented anatase structure, with an average crystallite size of a few tens of nanometers. Both layers present a very low porosity (near 3%), as environmental ellipsometric porosimetry (EEP) measurements demonstrate (Figure S1). Crystallization of the originally amorphous TiO_2 layer gives rise to a small amount of homogeneously distributed micro-cracks. The bimodal pore structure of the middle mesostructured layer can be readily appreciated in Figure 2b. The effect of incorporating PEG and THF into the precursor solution is to produce large, interconnected mesopores (75 nm average diameter) with local ordering, while the smaller pores (15 nm average diameter) present longer range ordering. Most importantly for the goal herein pursued, the interconnecting windows between pores present an average size of ≈ 10 nm, significantly larger than the size of the CSQDs employed, whose hydrodynamic diameter nominal value is 8.1 nm including the organic capping. (See supporting information, Figure S2) Details of the EEP analysis and the actual sorption-desorption curves are provided in the experimental and the supporting information sections (Figure S3). Consistently, multilayers in which disordered mesostructured mid layers were built up showed no sharp peaks in the low-angle diffractograms, in agreement with HRTEM images (Figure 2b). The thickness of such mid-layer could be precisely tuned in the range comprised between 50 nm and 150 nm varying either the concentration of the precursor or the substrate withdrawal speed.

So-built mesostructured optical resonators were used as matrices to incorporate luminescent particles from a liquid dispersion.

Please notice that the integration of quantum dots *selectively*

within an optical cavity of the sort herein described cannot be performed during the deposition of the constituent films, since Q-dots are not stable versus the thermal treatments usually needed to endow mechanical robustness to the multilayer structure. An approach to incorporate quantum dots in photonic structures built by sol-gel was previously developed by Jasieniak et al.³² In that case, an asymmetric optical resonator was completed using a silver mirror to prevent the thermal treatment of the sample once the CSQDs were incorporated into an upper layer deposited onto a dense TiO_2 - SiO_2 Bragg mirror.

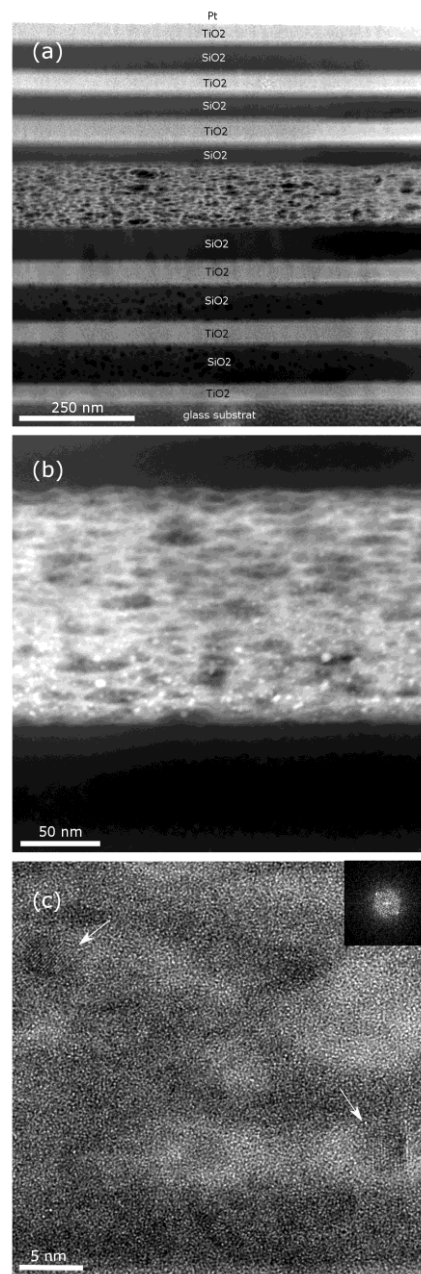


Fig. 2 Cross sectional ADF STEM images of the (a) complete 13-layers optical resonator (b) ADF image of the porous layer sandwiched in the middle of the optical structure with CSQD appearing as bright spots due to their high atomic number. (c) HRTEM image of one region of the porous layer. The arrows indicate regions, whose image Fourier transform is depicted in the inset, in which crystalline planes of a nanoparticle can be seen.



In our case, we take advantage of the porosity of the resonator to
infill the CSQDs after the multilayer is finalized. The thicknesses
of the constituent films were chosen so that the spectral position
of the optical resonance coincides with the emission band of the
CdSe/ZnS quantum dots used for this experiment. The luminescent
nanocrystals were introduced within the mesopores of the middle
layer by immersion. After soaking the bimodal pore size distribution
resonators overnight in a suspension of said nanoparticles, we
found consistent evidence by electron microscopy and photoluminescence
that diffusion and hence integration of CSQDs in the porous optical
cavity has occurred. In Figure 2.b, the location of clusters of CSQDs
within the mesostructured middle layer can be seen from the chemical
contrast in the ADF STEM image. The HRTEM image in Figure
2c reveals the crystalline structure of the CSQDs present in the
pores. An EDX linescan performed over such particles within the
mesostructured layer shows a clear cadmium signal at the particles
position (Figure S4).

The loss of Se signal may be attributed to the TEM sample
preparation, which might cause oxidation or degradation of the
delicate CSQDs. Nevertheless, TEM images prove that CSQD are
inside the porous optical cavity layer. For the sake of comparison,
we prepared mesostructured resonators using only F127 as
macromolecular template and subjected them to the same
infiltration process. In that case, no trace of luminescence was
detected from the resonators, which indicated that no luminescent
particles were incorporated, in good agreement with the
interconnecting window or neck size estimated from the
corresponding EEP analysis, which reveals sizes of 8.6nm and
4.6nm for pores and necks respectively. (Figure S5). So, even
when pore sizes in the resonator prepared using only F127 in the
precursor solution are large enough to host the CSQDs, pore
necks are smaller than the average emitter size and thus
infiltration is prevented. This dramatic effect has also been
observed in the adsorption of bulky enzymes in the whole pore
volume of this type of films,^{33,34} demonstrating one of the
advantages of the use of co-templates towards larger mesopore
systems with improved interconnectivity.

Regarding the optical characterization of the resonator, both the
dense multilayers and the mesostructured middle layer were
designed to yield an optical cavity mode at wavelengths matching
the photoemission band of the CdSe/ZnS nanocrystals infiltrated
in the structure, so the effect of the interplay between both
phenomena could be studied. Reflectance maxima result from the
constructive interference of beams of light in a certain
wavelength range reflected at the different interfaces existing
between each pair of layers in a periodic structure. The position
of this Bragg peak depends on the refractive indices and
thicknesses of the layers. An intermediate layer of different
optical thickness breaks the symmetry of the system and is
capable of hosting resonant modes, which can be recognized by
the opening of a transmission window at photonic band gap
frequencies. In this way, the whole structure behaves as an optical
cavity, whose transmittance at normal incidence is plotted in
Figure 3. From the analysis of the transmittance peak at around
 $\lambda=540$ nm, we estimate a cavity quality factor of $\lambda/\Delta\lambda \approx 30$.
Fitting of these spectra using a MATLAB code,³⁵ allowed us
estimating the thickness of each constituent layer, the unit cell of

the dense Bragg mirror being composed of 110 ± 8 nm thick SiO₂
and 58 ± 6 nm thick TiO₂ layers. The mesostructured middle layer
has a thickness of 180 ± 15 nm and an average refractive index of
 $n=1.70$ in the visible range. The wavelength dependent refractive
index and extinction coefficient curves employed for this fitting
are provided as supporting information (Figure S7). We use this
data to calculate the spatial distribution of the squared magnitude
of the electric field within the stack when a plane wave impinges
in an arbitrary direction onto the multilayer. In Figure 3b we plot
the particular $|E|^2$ attained for normal incidence. Different optical
resonances can be identified as bright spots within the multilayer.
Specifically, field intensity reinforcement is observed at spatial
and spectral positions coincident with the photonic band gap
edges (approximately at $\lambda=435$ nm and $\lambda=665$ nm) and the
optical cavity mode (ca. $\lambda=540$ nm).

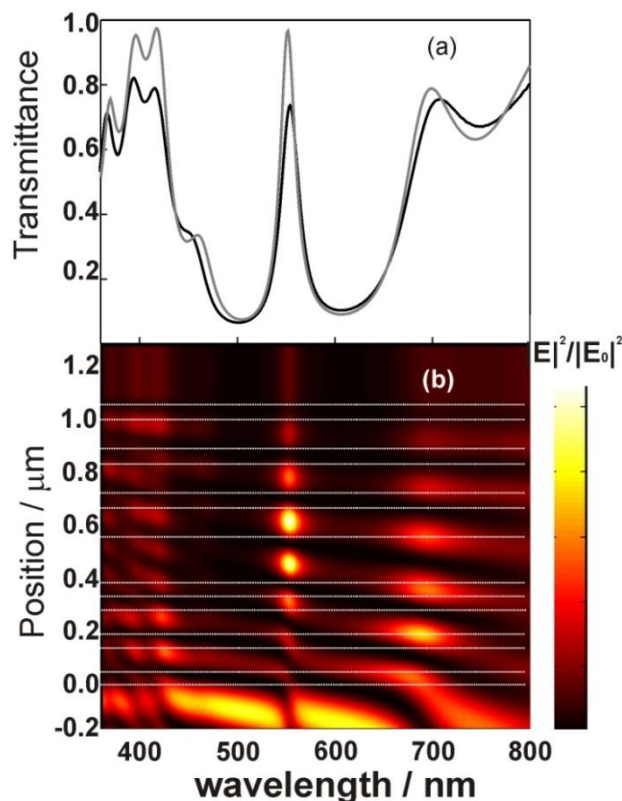


Fig. 3 (a) Experimental (black solid line) and theoretical (gray solid line) transmittance of the porous optical resonator. (b) Calculated spatial (vertical axis) and spectral (horizontal axis) distribution of the normalized electric field intensity enhancement ($|E|^2/|E_0|^2$) along a cross section of the mesoporous optical resonator. Calculations were carried out considering the multilayer parameters extracted from the fittings to the transmittance measurement. Horizontal dashed lines are guides to the eye to delimit the interfaces between layers in the optical resonator. The bottom of the images is the incoming medium (air) and the top, is the glass side. Incidence is 0 degrees.

Photoluminescence (PL) spectra were collected at different exit
angles from the CSQD loaded optical resonator using an
excitation wavelength of 370 nm, which falls outside the spectral
range for which resonant modes are expected, impinging along a
fixed incident direction. It should be noted that similar PL spectra
were taken from spots located at different positions on the
resonator surface, including the central part of the film, hence
further confirming that the density of CSQDs within the porous



optical cavity is homogeneous as a result of the good lateral diffusion of the particles throughout the middle layer. On the grounds of Lorenz reciprocity principle,^{36,37} the PL radiated from an emitter within the multilayer is expected to be enhanced along those directions and wavelengths for which a plane wave hypothetically traveling from the detector to the source would show a high field intensity at the position where the emitter is located. Analogously, the emission will be depleted at those wavelength ranges and along those extraction directions for which the corresponding field distribution display a dark region where the emitter is located.

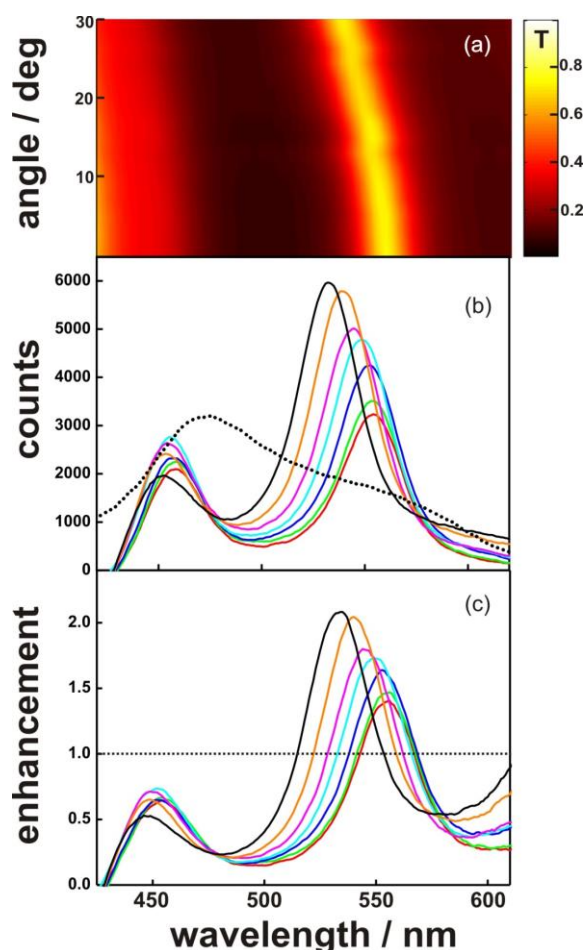


Fig. 4 (a) Experimental ballistic transmittance spectra of the porous resonator collected at different angles. Measurements were obtained each 2 degrees. Color map picture is built interpolating two consecutive measurements. (b) Luminescence spectra of the porous optical structure collected at 0 (black line), 15 (red line), 20 (blue line), 25 (green line) and 30 (brown line) degrees. Luminescence at 20 degrees obtained for a reference sample is plotted in dashed line. (c) Spectral enhancement of the luminescence at different angles (color code is the same described in Figure 4b)

These effects are actually observed in our measurements and are displayed in Figure 4. In Figure 4a we plot the experimental transmittance of a CSQD loaded resonator *versus* angle of incidence with respect to the multilayer surface normal. As wavelengths at which hot spots appear in the field distribution pattern coincide with those of transmission maxima at each specific angle, PL enhancement is expected to blue-shift with the

angle of collection with respect to surface normal, just like the cavity mode does in Figure 4a. Figures 4b shows the actual PL spectra at different angles of collection and the ratio between the PL of the resonator and that of an unstructured reference sample, respectively. This consists of a CSQD embedded mesostructured film, similar to that sandwiched between the two multilayers in the resonator, deposited on a glass substrate.

Emission reinforcement and depletion spectral regions match bright and dark regions in the transmittance plot of Figure 4a, as expected. For the sake of comparison, the emission of the CSQDs from the isolated mesostructured thin film (dashed line in Figure 4. b) remains approximately constant with the collection angle (Figure S6 in the supporting information). At the same time, PL quantum yield (ratio between the total number of emitted and absorbed photons) of the resonator embedded with CSQDs do not show any significant difference with that of the bare nanocrystals. This behaviour is the one expected for a planar structure for which no lateral confinement has been built up and hence for which no significant modifications of the Purcell factor are expected. In this case, channelling of emitted light along those directions for which there are more available modes does not imply that the intensity of emitted light integrated over all possible directions is larger. Please notice that the reciprocity principle can only provide a qualitative description of the PL enhancement and depletion phenomena observed, and that, in order to fully describe the effect of the inclusion of the CSQDs within the optical cavity on their emission properties, a detailed calculation of the local density of states at each spot from which emission takes place would be required.³⁸ In this regard, as the LDOS is expected to have a non-constant spectral and spatial profile within the resonator, the recorded PL spectra are inhomogeneous since they are the result of collecting light emitted from spots with different photonic environment.

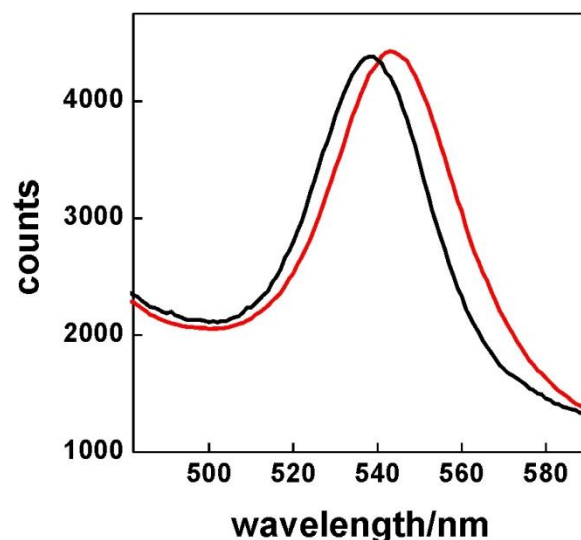


Fig. 5 Luminescence spectra of the porous optical structure collected at 0 degrees without (black line) and with (red line) ethanol infiltrated in the structure

Finally, we confirmed that the responsive character of the porous structure against changes in the environment was maintained after QD infiltration, revealing that the void network remained open and accessible. In Figure 5 we plot the photoluminescence spectra



of the CSQD embedded optical resonator when exposed to liquid ethanol (red line). The 6 nm red-shift observed in the emission peak position is the result of the interplay with the spectral shift in the resonant wavelength of the optical cavity, which is in turn due to the increase of the refractive index of the middle layer upon filling with the ethanol flowing from the lateral sides of the film.

Experimental Section

Preparation of the porous mesostructured optical resonators.

In order to prepare porous optical resonators, first a periodic structure made of alternate layers of TiO_2 and SiO_2 was deposited by dip-coating from sol-gel precursor suspensions. Then, this initial structure was coated with a mesostructured layer by dip-coating a supramolecularly templated precursor dispersion. Finally, a second periodic structure was grown onto the mesoporous layer.

Synthesis and deposition of dense TiO_2 and SiO_2 layers

TiO_2 and SiO_2 sols were prepared after procedures reported elsewhere.^{24,25} Glass slides cleaned with a HNO_3 solution (5%) and rinsed with water and ethanol were used. For the synthesis of the SiO_2 sol, we used 4.5 ml of tetraethylorthosilicate (TEOS, Merck) and 34 ml of absolute ethanol (EtOH). After some minutes under vigorous stirring, 1.72 ml of H_2O and 0.08 ml of HCl 0.05N were added. To prepare the TiO_2 sol, 2.5 ml of titanium isopropoxide (TIPT, 97%, Aldrich), 0.156 ml of HNO_3 (1M), 78 μl of H_2O and 25 ml of isopropyl alcohol (iPrOH, Aldrich) were mixed. The dip-coater employed to deposit the multilayer from these precursors was a HWTL-01-A, from MTI Corporation. The withdrawal speed of the substrates, from the suspension it was immersed in, was varied between 0.5 and 2 mm/s. This permitted us to obtain thicknesses in the range between 50 nm and 200 nm. Stabilization temperature was fixed at 300°C for 1h, and each layer in the stack was subjected to this treatment before a new one was deposited on top of it.

Synthesis and deposition of mesoporous layers

The porous layer precursor sol was synthesized using a novel recipe that allows to obtain bigger pore size using a low-cost polymer. A quantity of 12.6 grams of a TiCl_4 :BuOH solution (molar ratio $\text{Ti}/\text{BuOH}=1:40$) was used to dissolve 0.28 g of F127. After 10 minutes of vigorous mixing, 0.72g of H_2O were added and then 0.24 g of polypropylene glycol 4000 (Alfa Aesar GMBH&Co KG) was incorporated to the sol. Finally, 4 ml of tetrahydrofuran (THF) is mixed with the sol during 30 minutes. This precursor suspension was dip-coated onto the first dense multilayer, preserved at 50% relative humidity during 24h and then sequentially heated at 60°C (12h), 120°C (12h) and 200°C (2h). Once the mesostructure was stabilized, the second multilayer was deposited onto it following a similar method than the one employed to prepare the first one. The whole structure was thermally treated at 350°C for 1h (ramp 1°C/min) to remove the supramolecular compound used as template.

Porous mesostructured resonators were soaked in a toluene solution of CdSe/ZnS quantum dots (LumidotTM480, Aldrich) during 24 hours. After that, the samples were rinsed with fresh

toluene to eliminate excess of quantum-dots. Monolayers of same thickness as the middle layer used to build the optical cavity and possessing a similar porous mesostructure were treated in the same way in order to obtain a luminescent film whose emission can be compared to that of the resonator.

Structural characterization

Pore and neck size of individual films were characterized through environmental ellipsometric porosimetry (EEP). A SOPRA GES5A apparatus was used to attain water adsorption-desorption isotherms at 298 K by EEP. From the fitting of the ellipsometric parameters in the visible range (400-800 nm) we estimated film thickness and refractive index (real part). The film refractive index was described according to a modified Cauchy equation. The WinElli 2 software (SopraInc) is employed to convert variations of n with P/P° into filled pore volume. Kelvin equation is used to extract pore and neck size distributions.²⁹

For TEM analysis, a Lamella was cut with a FEI Helios Nanolab 400s with a subsequent milling in a Fischione Nanomill to remove destroyed surface layers. TEM and STEM experiments were conducted at FEI Tecnai F20 (ER-C) operated at 200 kV.

Analysis of the optical properties and their environmental dependence.

Transmittance spectra were obtained using an UV-VIS scanning spectrophotometer (UV-2101PC, Shimadzu). Transmission angle measurements were acquired supporting the sample onto a rotating stage (part number M55-028, Edmund Optics). Excitation and emission spectra were collected using a spectrofluorometer (Fluorolog-3, Horiba JobinYvon) with a Xe450 W lamp source. PL experiments were made collecting light in front face configuration with a fixed angle of 22.5° between the entrance and the exit slits. Sample was placed in the same rotating stage described above to obtain the emission at different angles. Ethanol-soaking experiments were made by placing a drop of ethanol onto the sample and pressing it gently with a cover slide to force infiltration.

Samples were constructed on a 75 mm x 25 mm low fluorescence glass. Metal oxide precursors liquid dispersions used to fabricate dense and porous layers cover an 80% of that area. Then samples were cut in three equal parts (regions near to the borders of the glass were discarded) and immersed in q-dots toluene dispersions. Final size of the samples is 15mm x 15mm. In all the cases, different series of photoluminescence measurements were taken from the side to the center of the sample to probe the degree of penetration of the SCQDs throughout the middle porous layer. We find that luminescence is not dependent of the region of the sample measured, thus proving the homogeneous lateral infiltration of the resonator.

Conclusions

We have demonstrated a synthetic route to selectively integrate colloidal semiconductor quantum dots in a porous optical cavity surrounded by dense photonics crystals. This is, to the best of our knowledge, the first time that precise spatial confinement of colloidal quantum dots in a resonator is achieved, as their incorporation in this configuration has been prevented before due to the low chemical and thermal stability of the organically



capped nanocrystals. In order to do so, the mesostructure of the resonator was carefully devised to permit the efficient lateral infiltration of the nanocrystals by soaking. We demonstrate that the photoemission spectrum of the confined quantum dots is determined by the cavity resonance, which opens the door to precise tailoring of the optical properties of this type of nanocrystals by confinement in porous matrices in which optical resonances are built up. We foresee our work may ease the integration of colloidal semiconductor quantum dots into new light emitting coatings with improved spectral and directional properties.

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Notes and references

^a Instituto de Ciencia de Materiales de Sevilla, Consejo Superior de Investigaciones Científicas-Universidad de Sevilla, Américo Vespucio 49, 41092 Sevilla, Spain

^b Institute of Energy and Climate Research: Fundamental Electrochemistry (IEK-9), Forschungszentrum Jülich GmbH, D-52425 Jülich, Germany Address here.

^c Ernst Ruska-Centre for Microscopy and Spectroscopy with Electrons, Forschungszentrum Jülich GmbH, D-52425 Jülich, Germany

^d Gerencia Química – Centro Atómico Constituyentes, Comisión Nacional de Energía Atómica, Av. Gral Paz 1499 (B1650KNA) San Martín, Buenos Aires, Argentina

† Footnotes should appear here. These might include comments relevant to but not central to the matter under discussion, limited experimental and spectral data, and crystallographic data.

Electronic Supplementary Information (ESI) available: [details of any supplementary information available should be included here]. See DOI: 10.1039/b000000x/

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